Finite difference methods

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Introduction

Finite difference methods (FDM) are numerical methods for solving (partial) differential equations, where (partial) derivatives are approximated by finite differences.

Example

For $0 < |h| \ll 1$, we have the approximations

(i)
$$f'(x) \approx \frac{f(x+h) - f(x)}{h}$$
,
(ii) $f'(x) \approx \frac{f(x) - f(x-h)}{h}$,
(iii) $f'(x) \approx \frac{f(x+h) - f(x-h)}{2h}$, and
(iv) $f''(x) \approx \frac{f(x+h) - 2f(x) + f(x-h)}{h^2}$

Let $a = x_0 < x_1 < \cdots < x_n = b$ be a uniform partition of the interval [a, b], where $x_k = a + hk$, $k \in \{0, 1, \dots, n\}$, and h = (b - a)/n.

The formulae on the previous page can be used to derive discrete approximations for

(i)
$$[f'(x_0), f'(x_1), \dots, f'(x_{n-1})]^{\mathrm{T}} = A\mathbf{f},$$

(ii) $[f'(x_1), f'(x_2), \dots, f'(x_n)]^{\mathrm{T}} = B\mathbf{f},$
(iii) $[f'(x_1), f'(x_2), \dots, f'(x_{n-1})]^{\mathrm{T}} = C\mathbf{f},$ and
(iv) $[f''(x_1), f''(x_2), \dots, f''(x_{n-1})]^{\mathrm{T}} = D\mathbf{f}$
using the point values $\mathbf{f} = [f(x_0), f(x_1), \dots, f(x_n)]^{\mathrm{T}}.$

Proof. Homework! :)

Note: By using a suitable quadrature rule (e.g., the left-, mid-, or right-point rule or the trapezoidal rule), one can obtain a discretized approximation for the antiderivative as well!

Finite difference methods for the equation y'(t) = f(t, y(t))

Let's consider the IVP

$$\begin{cases} y'(t) = f(t, y(t)), & t \ge 0, \\ y(0) = y_0. \end{cases}$$

To ensure (local) unique solvability, it is assumed that f is continuous and that $f(t, \cdot)$ is Lipschitz-continuous [Picard–Lindelöf theorem].

How do we get rid of the derivative? By integration:

$$y(t) = y_0 + \int_0^t f(\tau, y(\tau)) \,\mathrm{d}\tau.$$

Bad news: We have just turned our differential equation into an integral equation, which we still need to solve for y(t).

Good news: We know how to do numerical integration!

We are interested in solving the IVP for $t \in [0, T]$, T > 0. Let $0 = t_0 < t_1 < \cdots < t_n = T$ be an equispaced partition of the interval [0, T], $t_k = hk$, $k \in \{0, 1, \ldots, n\}$, and h = T/n.

Let us denote $y(t_k) = y_k$. We can recast the integral equation on the previous page as

$$y_{k+1} = y_k + \int_{t_k}^{t_{k+1}} f(t, y(t)) dt$$
 for $k = 0, 1, ..., n-1$.

Let's try different quadratures for approximating the integral above!

Remark. We omit the detailed error analysis of the resulting formulae and focus on some heuristics of the solution methods.

One-point Riemann sum approximations over $[t_k, t_{k+1}]$

Left-point rule:

$$y_{k+1} \approx y_k + hf(t_k, y_k).$$
 (Euler's method)

Right-point rule:

$$y_{k+1} \approx y_k + hf(t_{k+1}, y_{k+1}).$$
 (Implicit Euler's method)

Midpoint rule:

$$y_{k+1} pprox y_k + hf(t_k + rac{h}{2}, y_k + rac{h}{2}f(t_k, y_k)).$$
 (Runge–Kutta method)

$$y_{k+1} \approx y_k + \frac{h}{2}(f(t_k, y_k) + f(t_{k+1}, y_{k+1}))$$

 $pprox y_k + \frac{h}{2}f(t_k, y_k) + \frac{h}{2}f(t_k + h, y_k + hf(t_k, y_k)).$

$$y_{k+1} \approx y_k + \frac{h}{2}(f(t_k, y_k) + f(t_{k+1}, y_{k+1}))$$

$$\approx y_k + \frac{h}{2}f(t_k, y_k) + \frac{h}{2}f(t_k + h, y_k + hf(t_k, y_k)).$$

$$k_{1} = hf(t_{k}, y_{k}),$$

$$y_{k+1} \approx y_{k} + \frac{k_{1}}{2} + \frac{h}{2}f(t_{k} + h, y_{k} + k_{1}).$$

$$k_{1} = hf(t_{k}, y_{k}),$$

$$y_{k+1} \approx y_{k} + \frac{k_{1}}{2} + \frac{h}{2}f(t_{k} + h, y_{k} + k_{1}).$$

$$k_1 = hf(t_k, y_k), \ k_2 = hf(t_k + h, y_k + k_1),$$
$$y_{k+1} \approx y_k + \frac{k_1}{2} + \frac{k_2}{2}.$$

This is the second order Runge–Kutta method (RK2).

There is a similar relationship between the famous fourth order Runge–Kutta method (RK4) and Simpson's rule:

$$k_{1} = hf(t_{k}, y_{k}), \quad k_{2} = hf(t_{k} + \frac{h}{2}, y_{k} + \frac{1}{2}k_{1}),$$

$$k_{3} = hf(t_{k} + \frac{h}{2}, y_{k} + \frac{1}{2}k_{2}), \quad k_{4} = hf(t_{k+1}, y_{k} + k_{3}),$$

$$y_{k+1} \approx y_{k} + \frac{1}{6}(k_{1} + 2k_{2} + 2k_{3} + k_{4}).$$

The RK4 rule can be derived by first applying Simpson's rule and then balancing the "midstep" (here $y_{k+1/2} = y(t_{k+1/2})$ and $t_{k+1/2} = t_k + h/2$)

$$4f(t_{k+1/2}, y_{k+1/2}) = 2f(t_{k+1/2}, y_{k+1/2}) + 2f(t_{k+1/2}, y_{k+1/2})$$

by using two different methods:

$$\begin{aligned} f(t_{k+1/2}, y_{k+1/2}) &\approx f(t_k + \frac{h}{2}, y_k + \frac{1}{2}k_1), & (\text{left-point rule}) \\ f(t_{k+1/2}, y_{k+1/2}) &\approx f(t_k + \frac{h}{2}, y_k + \frac{1}{2}k_2). & (\text{right-point rule}) \end{aligned}$$

In summary

There exists a wide range of literature on FD methods for first order IVPs – and the theory extends nicely to first order systems of ODEs as well. The examples on pg. 7 highlight a couple of canonical cases:

- (A) Explicit methods,
- (B) Implicit methods.

While simple and intuitive, explicit methods are never unconditionally stable. For example, when working with [stiff equations], certain implicit methods are able to capture the asymptotics of the solution regardless of the step size.

Another class of FD methods is represented by Runge–Kutta methods, where the FD solution is adjusted based on point evaluations at additional intermediate points between the discretization mesh.

For an introduction into the approximation and stability theory of first order FD methods, I highly recommend the monograph [Griffiths and Higham]!

Finite difference methods for Poisson's equation $-\Delta u = f$

Let

$$\Delta u(x,y) = \frac{\partial^2}{\partial x^2} u(x,y) + \frac{\partial^2}{\partial y^2} u(x,y)$$

denote the Laplacian.

In the following, we discuss the solution of Poisson's equation

$$-\Delta u = f$$
 in a bounded domain $\Omega \subset \mathbb{R}^2$

with the Dirichlet boundary condition

$$u = g$$
 on $\partial \Omega$.

where f and g are given functions.

We need to discretize Poisson's equation in order to find a suitable matrix approximation $A\mathbf{u} = \mathbf{m}$, which we can then solve numerically.

We can obtain an approximate formula for Δu near (x, y) by developing Taylor expansions of u(x, y) with respect to x and y variables, respectively. Assuming sufficient smoothness, we obtain

$$u(x + h, y) = u(x, y) + h\partial_x u(x, y) + \frac{1}{2}h^2\partial_x^2 u(x, y) + \frac{1}{6}h^3\partial_x^3 u(x, y) + \mathcal{O}(h^4),$$

$$u(x - h, y) = u(x, y) - h\partial_x u(x, y) + \frac{1}{2}h^2\partial_x^2 u(x, y) - \frac{1}{6}h^3\partial_x^3 u(x, y) + \mathcal{O}(h^4).$$

The sum of these is

$$u(x+h,y)+u(x-h,y)=2u(x,y)+h^2\partial_x^2u(x,y)+\mathcal{O}(h^4)$$

Developing the Taylor expansion w.r.t. the x variable yielded

$$u(x+h,y) + u(x-h,y) = 2u(x,y) + h^2 \partial_x^2 u(x,y) + \mathcal{O}(h^4).$$
(1)

Proceeding analogously with the other component:

$$u(x, y + h) = u(x, y) + h\partial_{y}u(x, y) + \frac{1}{2}h^{2}\partial_{y}^{2}u(x, y) + \frac{1}{6}h^{3}\partial_{y}^{3}u(x, y) + \mathcal{O}(h^{4}),$$

$$u(x, y - h) = u(x, y) - h\partial_{y}u(x, y) + \frac{1}{2}h^{2}\partial_{y}^{2}u(x, y) - \frac{1}{6}h^{3}\partial_{y}^{3}u(x, y) + \mathcal{O}(h^{4})$$

$$\Rightarrow u(x, y + h) + u(x, y - h) = 2u(x, y) + h^{2}\partial_{y}^{2}u(x, y) + \mathcal{O}(h^{4}).$$
 (2)

By summing (1) and (2) together, we obtain

$$u(x+h,y) + u(x-h,y) + u(x,y+h) + u(x,y-h) = 4u(x,y) + h^{2}\Delta u(x,y) + \mathcal{O}(h^{4})$$

$$\Rightarrow \boxed{\Delta u(x,y) = \frac{u(x+h,y) + u(x-h,y) + u(x,y+h) + u(x,y-h) - 4u(x,y)}{h^{2}} + \mathcal{O}(h^{2}).}$$

Five-point stencil (quincunx):



$$\Delta u(x,y) \approx \frac{u(x+h,y) + u(x-h,y) + u(x,y+h) + u(x,y-h) - 4u(x,y)}{h^2}$$

Recovering missing pixels

	100	100	100	100	
20	$\mathbf{x}_{_{1}}$	$\mathbf{x}_{_{4}}$	\mathbf{x}_{γ}	x ₁₀	20
20	x ₂	X ₅	x ₈	X ₁₁	20
20	х ₃	Х ₆	х ₉	X ₁₂	20
	100	100	100	100	

During the lecture, we discussed the example presented in Poisson_FD_v2.pdf (see course page).

Figure: Image by Samuli Siltanen.

For a given (polygonal) computational domain, we could attempt to create a quincunx mesh where, for each individual mesh element (x, y), the neighbouring coordinates $(x \pm h, y)$ and $(x, y \pm h)$ are either

- Other mesh elements,
- Boundary points.





Example

Let's solve the Dirichlet problem

$$\begin{cases} \Delta u = 0 & \text{ in } \Omega, \\ u|_{\partial\Omega} = g & \text{ on } \partial\Omega, \end{cases}$$

where the computational domain Ω is the L-shaped domain on the left, the discretization mesh is denoted by black dots, and the boundary values of g are given in red at the discretization nodes lying on the boundary.

Let's give a numbering to the unknowns!



Finite difference methods





















$$\frac{1}{h^2} A \begin{bmatrix} u(\mathbf{x}_1) \\ \vdots \\ u(\mathbf{x}_{16}) \end{bmatrix} = \frac{1}{h^2} m$$







$$\frac{1}{h^2}A\begin{bmatrix}u(\mathbf{x}_1)\\\vdots\\u(\mathbf{x}_{16})\end{bmatrix}=\frac{1}{h^2}m$$



$$\frac{1}{h^2}A\begin{bmatrix}u(\mathbf{x}_1)\\\vdots\\u(\mathbf{x}_{16})\end{bmatrix}=\frac{1}{h^2}m$$



$$\frac{1}{h^2} A \begin{bmatrix} u(\mathbf{x}_1) \\ \vdots \\ u(\mathbf{x}_{16}) \end{bmatrix} = \frac{1}{h^2} m$$



$$\frac{1}{h^2}A\begin{bmatrix}u(\mathbf{x}_1)\\\vdots\\u(\mathbf{x}_{16})\end{bmatrix}=\frac{1}{h^2}m$$



$$\frac{1}{h^2}A\begin{bmatrix}u(\mathbf{x}_1)\\\vdots\\u(\mathbf{x}_{16})\end{bmatrix}=\frac{1}{h^2}m$$

Numerical solution:

>> A\m

should give

 $[u(x_1),\ldots,u(x_{16})]^{\mathrm{T}}=[-4,-3,-3,-2,-2,-1,-1,0,1,2,3,0,1,2,3,4]^{\mathrm{T}}.$

Compare this with the analytical solution u(x, y) = x - y. (They are exactly the same!)

Numerical experiment

Let's return to the isospectral drums discussed earlier during the course.



Figure: Isospectral drum shapes in 2D [Gordon, Webb, and Wolpert].

Two drums with clamped boundaries give the same sound if they have the same set of (Dirichlet) eigenvalues λ satisfying

$$-\Delta u = \lambda u$$
 in $D, \ u|_{\partial D} = 0, \quad D \subset \mathbb{R}^2$ bounded domain.

The quincunx pattern has extremely important applications in statistics as well!

https://www.youtube.com/watch?v=AUSKTk9ENzg

Bibliography

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